

MATHEMATICAL MODELS FOR CELL MOTILITY WITH NONLOCAL REPULSION FROM SATURATED AREAS.

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Cell motility importance

Cellular movements are critical for a wide number of biological processes:

- embryogenesis,
- tissue formation,
- wound healing,
- defense against infection,
- control cancer metastasis and birth defects, among other.

Hence, the cellular motion has become an *increasingly interesting topic within biomedical research*.

It is well known that if cells get too close to each other, they *repulse* , while if they get too far away from each other, they *attract*.

On the other hand, they also exhibit certain *random movement*.

Saturation and repulsion prevent to occur high accumulations at narrow zones or single points.

Non-local operators

From a mathematical point of view, *nonlocal operators are crucial* to understand the immense majority of biological processes, by their *ability to taking into account the effect of the surrounding environment* to describe what happens at certain point, in contrast to local differential operators.

Non-local advection is a mechanism present in the mathematical modelling of a wide range of biological phenomena. Animals explore their surroundings to find prey, avoid predators or aggregate in colonies, herds or swarms.

This non-local sensing, which in animals is due mainly to using smell, hearing or sight, it also occurs at cellular level, through the extension of *long thin protrusions*.

Armstrong-Painter-Sherrat model

They consider one *single population of cells* moving (in 1d) responding to a natural random motion and also due to adhesive forces between the cells.

The forces acting on the cells define a *conservative system*, where no cell birth or death is supposed to occur:

$$\frac{\partial u}{\partial t}(x, t) = \frac{\partial J}{\partial x}(x, t),$$

where $u(x, t)$ is the density of cells at position x and time t .

J is the sum of the diffusive and adhesion fluxes, $J = J_d + J_a$.

For simplicity, the authors take the following diffusive flux

$$J_d(x, t) = \frac{\partial u}{\partial x}(x, t),$$

‘the results reported in this paper are the consequence of the adhesion term and could be obtained even without a diffusion term in the model’

Armstrong-Painter-Sherrat model

Regarding the *adhesive flux*, it is proposed to be *proportional to the density of the cells and the forces* between them are *inversely proportional to cell size*

$$J_a(x, t) = \frac{\nu}{R} u(x, t) F(x, t),$$

where $\nu > 0$ is a viscosity coefficient and $R > 0$ *the sensing radius*.

The total force is the nonlocal gradient

$$\begin{aligned} F(x, t) &= (\nabla_{NL} u)(x, t) := \mathcal{G}^+ u(x, t) - \mathcal{G}^- u(x, t) \\ &= \int_x^{x+R} u(y, t) \omega(y - x) dy - \int_{x-R}^x u(y, t) \omega(x - y) dy. \end{aligned}$$

After adimensionalization, they get the evolution equation

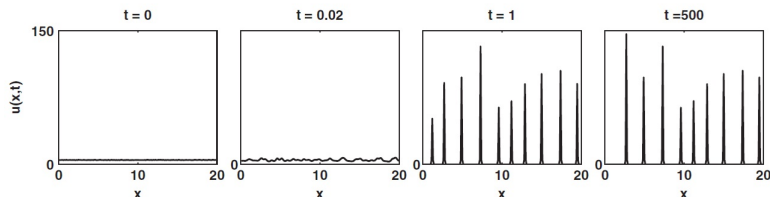
$$u_t(x, t) = u_{xx} - \gamma \frac{\partial}{\partial x} \left(u(x, t) (\nabla_{NL} u)(x, t) \right),$$

subject to *periodic boundary conditions*. Here $\gamma > 0$ *is the strength of the adhesion*

Armstrong-Painter-Sherrat model

Simulations of this model are performed, taking $R = 1$, $\omega(x) \equiv 1$ and $\gamma = 10$, on a domain of length 20 discretized into 200 mesh points.

The population of cells is *initially* assumed to be *uniformly distributed*, *perturbed by a small amount of noise*.



Over time the population develops a number of peaks in cell density. As time progresses, these *peaks are seen to coarsen*, becoming steeper and more widely spaced.

Review on APS model

Carrillo et al focus on APS model, to avoid high accumulations. With a *porous media* diffusion they consider

$$J_a(x, t) = (1 - u(x, t))u(x, t) \int_0^R \left(u(x + r, t) - u(x - r, t) \right) \omega(r) dr,$$

where the crowding capacity has been rescaled to $M = 1$, taking also the initial data below one.

In fact, the variable u of this model can be interpreted as the *volume fraction of cells*.

The factor $(1 - u(x, t))$ reduces their velocity as the area becomes gradually more crowded at x , producing *saturation effect*.

Further *repulsion effects* could be included through ω taking negative values, but depending on the distance between cells.

J. A. Carrillo, H. Murakawa, M. Sato, H. Togashi, & O. Trush, A population dynamics model of cell-cell adhesion incorporating population pressure and density saturation, J. Theoret. Biol. 2019

Questions about the previous models

- Is it feasible to include *repulsion effects, depending on the density of population as part of the drift term?*
- Is it possible to compute the *relative importance between diffusion and interaction processes?*
- How does the dynamics evolve in *confinement*?
There exist types of cells that are spontaneously motile only in confinement.

R.J. Hawkins, M. Piel, G. Faure-Andre, A. M. Lennon-Dumenil, J. F. Joanny, J. Prost, & R. Voituriez, Pushing off the walls: A mechanism of cell motility in confinement, Phys. Rev. Lett., 2009.

Description of a discrete model

Let $\{1, \dots, N\}$ be the cells.

At time $t = 0$ we assign a real number $x_i(0) \in \overline{\Omega} = [-L, L]$, representing the *initial position* of the cell $i \in \{1, \dots, N\}$.

Let us subdivide the interval $[-L, L]$ in $\{I_j\}, j \in \{1, \dots, M\}$, a family of intervals of length h .

Denote as

$$s_j(t) = \frac{\#\{i : x_i(t) \in I_j\}}{N}, \quad j \in \{1, \dots, M\},$$

the proportion of cells with position within the interval I_j , for $j \in \{1, \dots, M\}$.

Description of a discrete model

Fix the *characteristic time step* Δt . Then, for $j = 1, \dots, M$

$$s_j(t + \Delta t) = s_j(t) + \Delta t \left[Q_{\text{diff}}^h \left(G_{\text{diff}}(j, t) - L_{\text{diff}}(j, t) \right) + Q_{\text{int}}^h \left(G_{\text{int}}(j, t) - L_{\text{int}}(j, t) \right) \right],$$

- $G_{\text{diff}}(j, t), G_{\text{int}}(j, t)$ stand for probability *gain terms*,
- $L_{\text{diff}}(j, t), L_{\text{int}}(j, t)$ represent *loss terms*,

due both to diffusion and interaction effects, respectively.

- The coefficients Q_{diff}^h and Q_{int}^h account for the *frequency of diffusion and interaction processes*, respectively, per unit of time.
- *The dynamics is isolated.*

Random motion

Assuming that the jump of a cell is of length h , the proportion of cells arriving to I_j due to diffusion comes from adjacent intervals.

The transition probabilities to interval I_j are

$$P_{I_{j+1} \rightarrow I_j} = p/2 = P_{I_{j-1} \rightarrow I_j},$$

while the probability of remaining at I_j is $1 - p$. Namely,

$$G_{\text{diff}}(t, j) = s_{j+1}(t)P_{I_{j+1} \rightarrow I_j} + s_{j-1}(t)P_{I_{j-1} \rightarrow I_j} = \frac{p}{2}s_{j+1}(t) + \frac{p}{2}s_{j-1}(t).$$

On the other hand, the cells in I_j will leave it to travel to either I_{j-1} or I_{j+1} with respective probabilities

$$P_{I_j \rightarrow I_{j-1}} = p/2 = P_{I_j \rightarrow I_{j+1}}.$$

Thus, the proportion of cells leaving I_j is

$$L_{\text{diff}}(t, j) = s_j(t)(P_{I_j \rightarrow I_{j-1}} + P_{I_j \rightarrow I_{j+1}}) = p s_j(t).$$

Finding the underlying PDE

To this end, we introduce $u_h(x, t) : [-L, L] \times [0, \infty) \rightarrow \mathbb{R}_0^+$ such that

$$s_j(t) = \int_{I_j} u_h(x, t) dx = h u_h(x_j, t) + O(h^3),$$

being x_j the midpoint of each I_j .

In other words, u_h restricted to the interval I_j behaves as $s_j(t)/h$.

We rewrite the gain and loss in terms of the density u_h .

Assuming that

$$u_h \rightarrow u \text{ as } h \rightarrow 0 \text{ (in some sense)}$$

our aim is to deduce (formally) the *PDEs verified by u* .

Diffusive motion

We write the previous balance in terms of the density u_h ,

$$\begin{aligned}
 G_{\text{diff}}(t, j) - L_{\text{diff}}(t, j) &= \frac{p}{2} s_{j+1}(t) + \frac{p}{2} s_{j-1}(t) - p s_j(t) \\
 &= h \frac{p}{2} \left[(u_h(x_{j+1}, t) - u_h(x_j, t)) \right. \\
 &\quad \left. + (u_h(x_{j-1}, t) - u_h(x_j, t)) \right] \\
 &= h^3 \frac{p}{2} \left[\Delta_h u_h(x_j) \right],
 \end{aligned}$$

where

$$\Delta_h v(x) = \frac{v(x+h) - 2v(x) + v(x-h)}{h^2}$$

is the discrete Laplacian.

Aggregation through nonlocal gradient

The *proportion of cells arriving to l_j* are

$$G_{\text{int}}(t, j) = s_{j+1}(t)P_{l_{j+1} \rightarrow l_j}(t) + s_{j-1}(t)P_{l_{j-1} \rightarrow l_j}(t),$$

where the *transition probabilities* are given by

$$P_{l_{j+1} \rightarrow l_j}(t) = \sum_{i=j-r}^j s_i(t)w((j+1-i)h), \quad P_{l_{j-1} \rightarrow l_j}(t) = \sum_{i=j}^{j+r} s_i(t)w((i-(j-1))h).$$

On the other hand, *the proportion of cells leaving l_j* are

$$L_{\text{int}}(t, j) = s_j(t) \left[P_{l_j \rightarrow l_{j-1}}(t) + P_{l_j \rightarrow l_{j+1}}(t) \right],$$

being

$$P_{l_j \rightarrow l_{j-1}}(t) = \sum_{i=j-1-r}^{j-1} s_i(t)w((j-i)h) \quad P_{l_j \rightarrow l_{j+1}}(t) = \sum_{i=j+1}^{j+1+r} s_i(t)w((i-j)h).$$

The number $r = r(h) = \lceil R/h \rceil$.

Aggregation through nonlocal gradient

We bring together the following terms

$$\text{Left} := s_{j+1}(t)P_{I_{j+1} \rightarrow I_j}(t) - s_j(t)P_{I_j \rightarrow I_{j-1}},$$

$$\text{Right} := s_{j-1}(t)P_{I_{j-1} \rightarrow I_j}(t) - s_j(t)P_{I_j \rightarrow I_{j+1}}.$$

By the definition of u_h , one has

$$s_i(t)w(|j+1-i|h) = \int_{I_i} u_h(y, t)w(|x_{j+1}-y|)dy + O(h^3),$$

hence, from

$$P_{I_{j+1} \rightarrow I_j}(t) = \int_{\bigcup_{i=j-r}^j I_i} u_h(y, t)w(x_{j+1}-y)dy + O(h^2) = \mathcal{G}^- u_h(x_{j+1}, t) + O(h^2).$$

Similarly,

$$P_{I_{j-1} \rightarrow I_j}(t) = \int_{\bigcup_{i=j}^{j+r} I_i} u_h(y, t)w(y-x_{j-1})dy + O(h^2) = \mathcal{G}^+ u_h(x_{j-1}, t) + O(h^2).$$

Aggregation through nonlocal gradient

One arrives then to

$$\begin{aligned}
 \text{Left} &= \left[hu_h(x_{j+1}, t) \right] \left[\mathcal{G}^- u_h(x_{j+1}, t) + O(h^2) \right] \\
 &\quad - \left[hu_h(x_j, t) \right] \left[\mathcal{G}^- u_h(x_j, t) + O(h^2) \right] \\
 &= h^2 \left[\frac{u_h(x_{j+1}, t) \mathcal{G}^- u_h(x_{j+1}, t) - u_h(x_j, t) \mathcal{G}^- u_h(x_j, t)}{h} \right] + O(h^3)
 \end{aligned}$$

and

$$\text{Right} = -h^2 \left[\frac{u_h(x_j, t) \mathcal{G}^+ u_h(x_j, t) - u_h(x_{j-1}, t) \mathcal{G}^+ u_h(x_{j-1}, t)}{h} \right] + O(h^3)$$

Aggregation through nonlocal gradient

Summing up, we have shown that

$$G_{\text{int}}(t, j) - L_{\text{int}}(t, j) = -h^2 \left[\mathcal{T}_{\text{aps}}^h(x_j, t) + O(h) \right],$$

where

$$\mathcal{T}_{\text{aps}}^h(x_j, t) := \partial_h^- \left[u_h(x_j, t) \mathcal{G}^+ u_h(x_j, t) \right] - \partial_h^+ \left[u_h(x_j, t) \mathcal{G}^- u_h(x_j, t) \right].$$

Hence
$$\lim_{h \rightarrow 0} \mathcal{T}_{\text{aps}}^h(x_j, t) = \frac{\partial}{\partial x} \left(u(x, t) (\nabla_{NL} u)(x, t) \right).$$

where recall that

$$\begin{aligned} (\nabla_{NL} u)(x, t) &:= \mathcal{G}^+ u(x, t) - \mathcal{G}^- u(x, t) \\ &= \int_x^{x+R} u(y, t) \omega(y - x) dy - \int_{x-R}^x u(y, t) \omega(x - y) dy. \end{aligned}$$

Aggregation through nonlocal gradient

The master equation reads as

$$\left[h \frac{u_h(x_j, t + \Delta t) - u_h(x_j, t)}{\Delta t} \right] = Q_{\text{diff}}^h h^3 \mathcal{T}_{\text{diff}}^h(x_j, t) - Q_{\text{int}}^h h^2 [\mathcal{T}_{\text{aps}}^h(x_j, t) + O(h)],$$

where $\lim_{h \rightarrow 0} \mathcal{T}_{\text{diff}}^h(x_j, t) = \frac{p}{2} \Delta u(x, t).$

Aggregation through nonlocal gradient

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where $\lim_{h \rightarrow 0} \mathcal{T}_{\text{diff}}^h(x_j, t) = \frac{p}{2} \Delta u(x, t).$

To preserve all of the effects in the limit PDE, we find the **constrains**

$$\lim_{h \rightarrow 0} h^2 Q_{\text{diff}}^h = \alpha_{\text{diff}}, \quad \lim_{h \rightarrow 0} h Q_{\text{int}}^h = \alpha_{\text{int}}.$$

Aggregation through nonlocal gradient

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To preserve all of the effects in the limit PDE, we find the **constraints**

$$\lim_{h \rightarrow 0} h^2 Q_{\text{diff}}^h = \alpha_{\text{diff}}, \quad \lim_{h \rightarrow 0} h Q_{\text{int}}^h = \alpha_{\text{int}}.$$

This yields the following nonlocal PDE

$$u_t(x, t) = u_{xx}(x, t) - \gamma \frac{\partial}{\partial x} \left(u(x, t) (\nabla_{NL} u)(x, t) \right)$$

and **isolated** boundary conditions

$$\begin{aligned} u_x(L, t) + \gamma u(L, t) \int_0^R u(L - r, t) w(r) dr &= 0 \\ -u_x(-L, t) + \gamma u(-L, t) \int_0^R u(-L + r, t) w(r) dr &= 0. \end{aligned}$$

Lagrangian perspective

Again, consider a population of N cells randomly distributed in $[-L, L]$, at initial positions $x_i(0)$, $i = 1, \dots, N$.

Select an arbitrary particle at $x_i(0)$ and follow its trajectory, $x_i(t)$.

The following system of stochastic differential equations describes *how the position of each of the N cells evolves along time*:

$$dx_i(t) = \frac{1}{N} \sum_{j \neq i} \mathcal{F}(x_i(t), x_j(t), S_a(x_i)) dt + \sqrt{2} \varepsilon B_t^i, \quad i = 1, \dots, N, \quad t > 0.$$

Here, $\mathcal{F}(x_i(t), x_j(t), S_a(x_i)) = S_a(x_i) F(|x_i - x_j|)$ represents the action of the particle j interfering on the trajectory of the particle i .

The coefficient $S_a(x_i)$ models the saturation or repulsion mechanisms.

The terms B_t^i , $i = 1, \dots, N$ are independent Brownian motion.

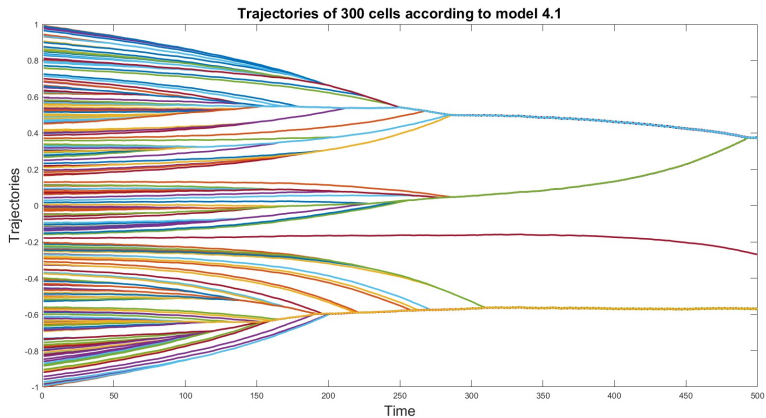
Lagrangian perspective

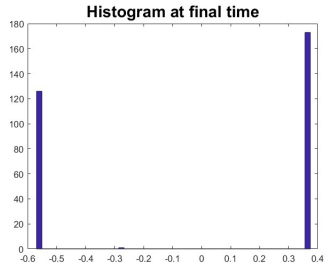
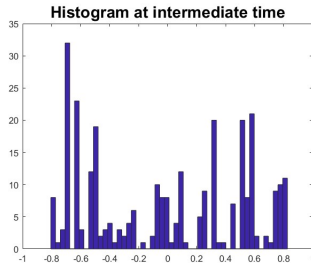
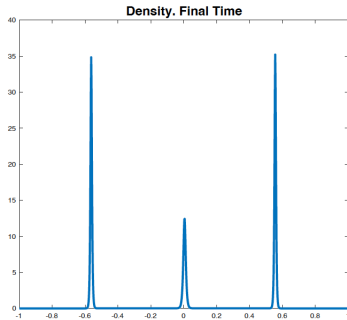
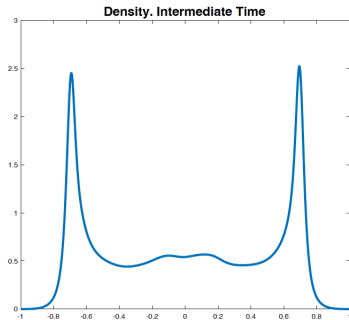
Let $\varphi : [-L, L] \rightarrow \mathbb{R}$ be an observable function and μ_t^N the empirical measure of the system of N cells:

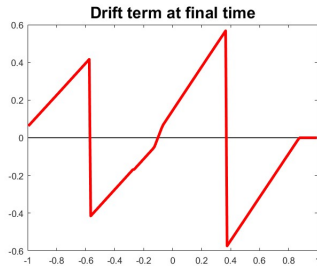
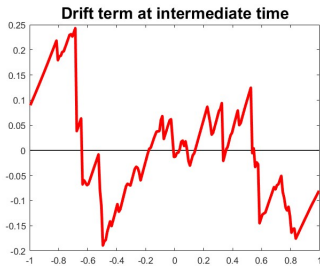
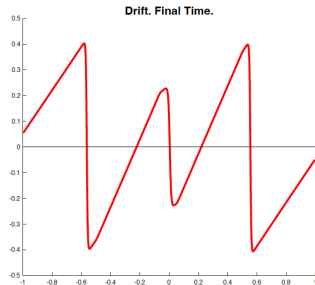
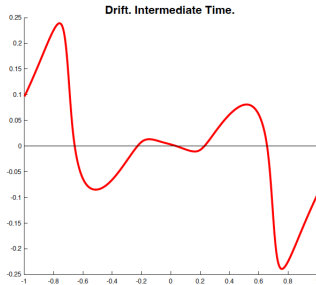
$$\mu_t^N = \frac{1}{N} \sum_{i=1}^N \delta_{x_i(t)},$$

where δ_x is the Dirac mass on $x \in \mathbb{R}$.

$$\begin{aligned} \langle \varphi, \delta_{x_i(t)} \rangle - \langle \varphi, \delta_{x_i(0)} \rangle &= \sqrt{2\varepsilon} \int_0^t \frac{d\varphi}{dx}(x_i(s)) dB_s^i \\ &+ \int_0^t \left[S_a(x_i(s)) \frac{1}{N} \sum_{j \neq i} F(x_i(s), x_j(s)) \frac{d\varphi}{dx}(x_i(s)) + \varepsilon^2 \frac{d^2\varphi}{dx^2}(x_i(s)) \right] ds. \end{aligned}$$







Lagrangian perspective: Local Saturation Model

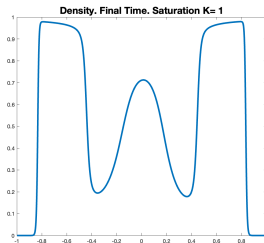
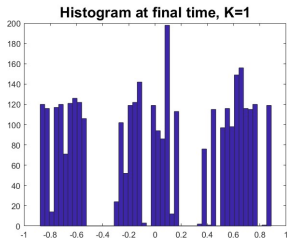
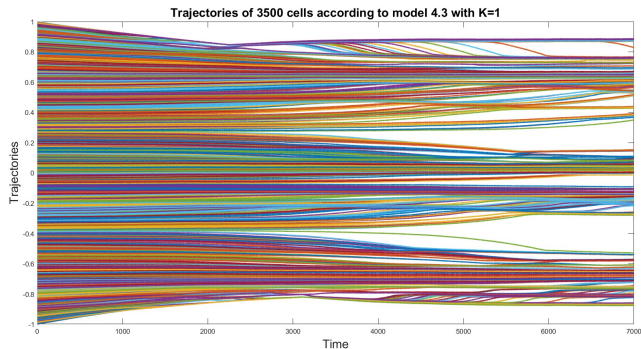
Take $a = a(N) \rightarrow 0$ as $N \rightarrow \infty$ and define

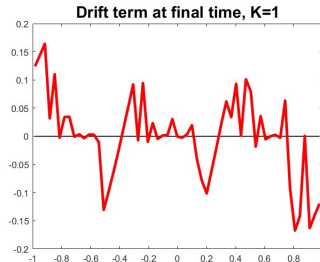
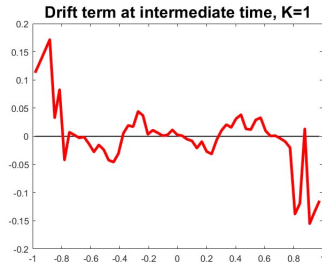
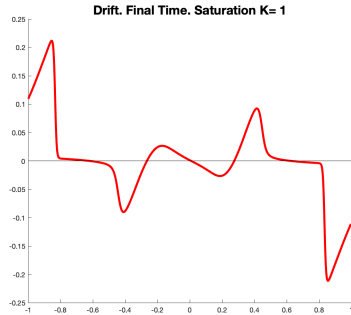
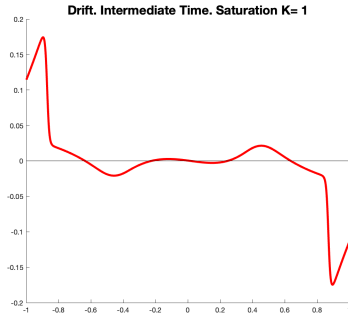
$$S_a(x_i(t)) = 1 - \frac{1}{2a} \langle \mathbb{1}_{|z-x_i(t)| < a}, \mu_t^N \rangle = 1 - \frac{1}{2a} \int_{x_i(t)-a}^{x_i(t)+a} d\mu_t^N(z).$$

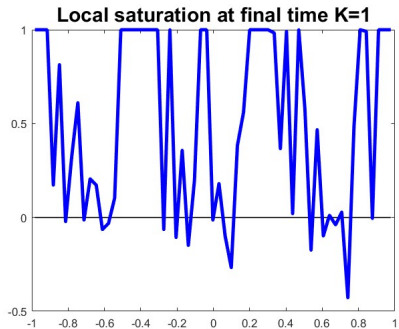
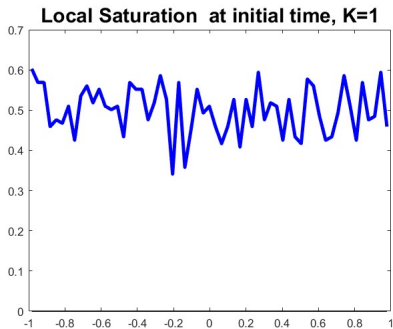
The measure μ_t^N is the empirical measure of the system of particles.

The coefficient $S_a(x_i(t))$ approaches to the local coefficient $1 - u$ as $N \rightarrow \infty$ in the PDE

$$u_t(x, t) = u_{xx}(x, t) - \gamma(1 - u(x, t)) \frac{\partial}{\partial x} \left(u(x, t) (\nabla_{NL} u)(x, t) \right)$$







A nonlocal aggregation/ repulsion coefficient

In the spirit of avoiding large accumulations, we propose to take a **nonlocal saturation coefficient** choosing $a = R$ the **sensing radius**. Namely,

$$\mathcal{F}(x_i, x_j, S_R(x_i)) = \left(1 - \frac{1}{2RK} \langle \mathbb{1}_{|z-x_i(t)| < R}, \mu_t^N \rangle \right) F(|x_i - x_j|),$$

for some value of the **crowding capacity** $0 < K \leq 1$.

Taking $N \rightarrow \infty$ it yields the PDE

$$u_t(x, t) = u_{xx}(x, t) - \gamma \left(1 - \frac{1}{K} \int_{x-R}^{x+R} u(y, t) \mathbb{1}_\Omega(y) dy \right) \frac{\partial}{\partial x} \left(u(x, t) (\nabla_{NL} u)(x, t) \right).$$

A nonlocal aggregation/ repulsion coefficient.

We analyze this new coefficient, which under the Euler perspective reads as

$$1 - \frac{1}{K} \sum_{i=j-l}^{j+l} s_i(t),$$

for some $K \in (0, 1]$.

It **can take negative values**, How could we then to define the transition probabilities among the intervals...?

A nonlocal aggregation/ repulsion coefficient.

$$\begin{aligned}
 P_{I_{j+1} \rightarrow I_j}(t) = & - \left(K - \sum_{i=j+1-l}^{j+1+l} s_i(t) \right)_{-} \sum_{i=j+2}^{j+2+r} s_i(t) w((i - (j+1))h) \text{ (repulsion term)} \\
 & + \left(K - \sum_{i=j+1-l}^{j+1+l} s_i(t) \right)_{+} \sum_{i=j-r}^j s_i(t) w((j+1-i)h) \text{ (aggregation term)}
 \end{aligned}$$

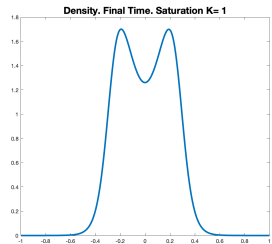
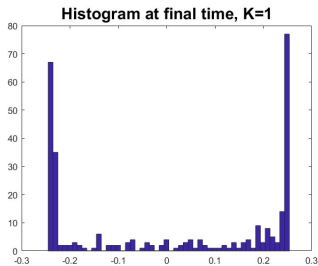
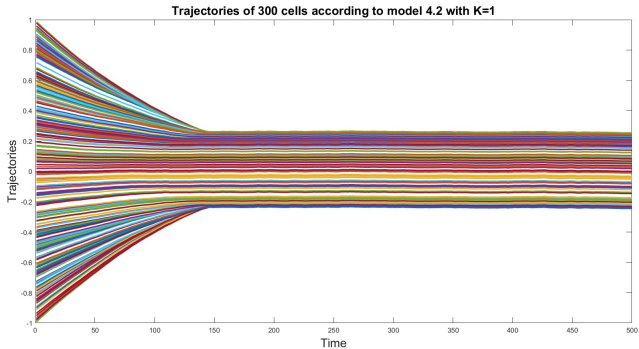
A nonlocal aggregation/ repulsion coefficient.

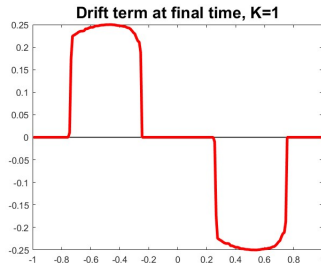
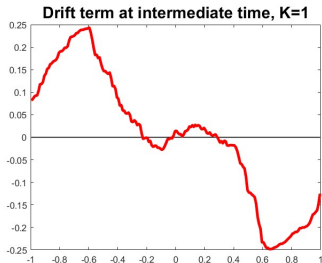
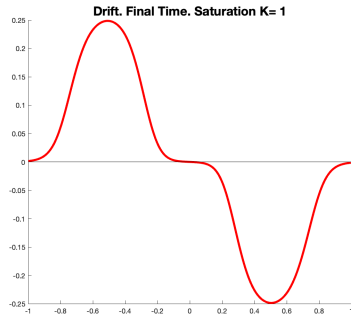
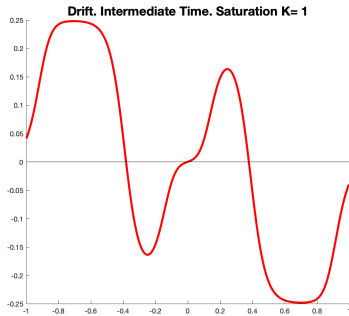
$$P_{I_{j+1} \rightarrow I_j}(t) = - \left(K - \sum_{i=j+1-l}^{j+1+l} s_i(t) \right)_{-} \sum_{i=j+2}^{j+2+r} s_i(t) w((i - (j+1))h) \text{ (repulsion term)}$$

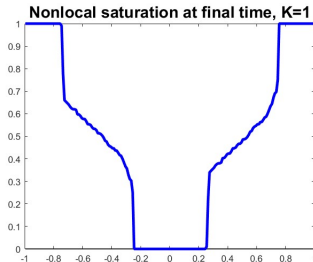
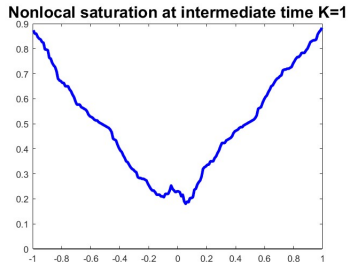
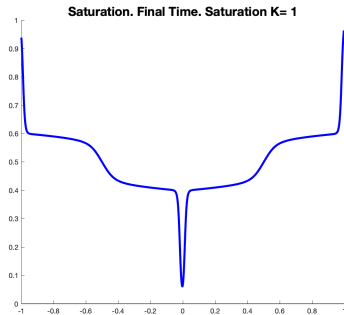
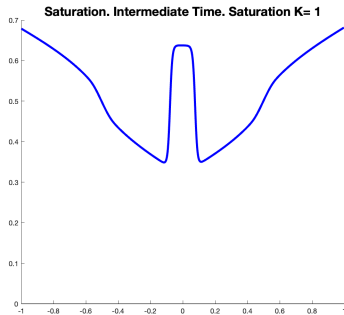
$$+ \left(K - \sum_{i=j+1-l}^{j+1+l} s_i(t) \right)_{+} \sum_{i=j-r}^j s_i(t) w((j+1 - i)h) \text{ (aggregation term)}$$

$$P_{I_{j-1} \rightarrow I_j}(t) = - \left(K - \sum_{i=j-1-l}^{j-1+l} s_i(t) \right)_{-} \sum_{i=j-2-r}^{j-2} s_i(t) w((j-1 - i)h) \text{ (repulsion term)}$$

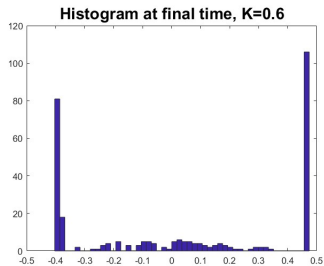
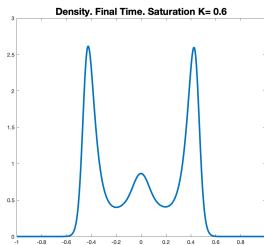
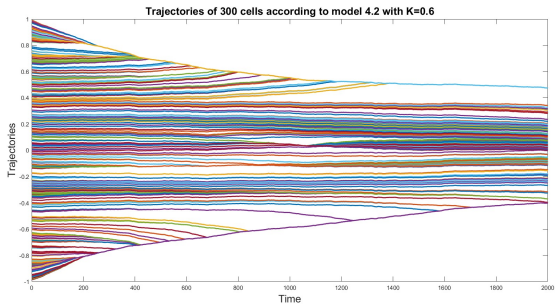
$$+ \left(K - \sum_{i=j-1-l}^{j-1+l} s_i(t) \right)_{+} \sum_{i=j}^{j+r} s_i(t) w((i - (j-1))h) \text{ (aggregation term)}$$

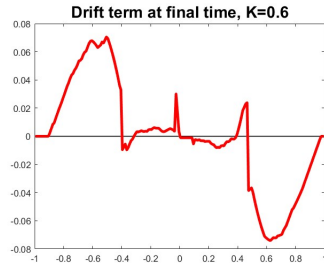
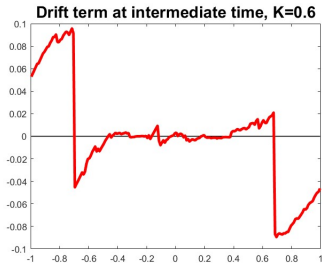
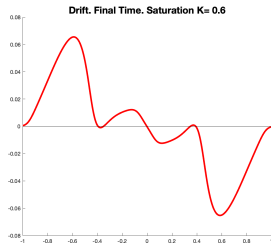
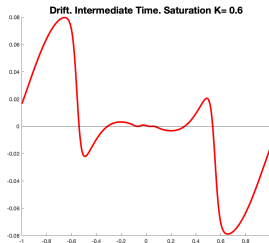


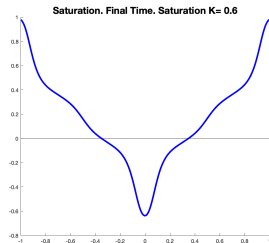
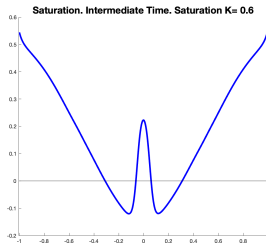




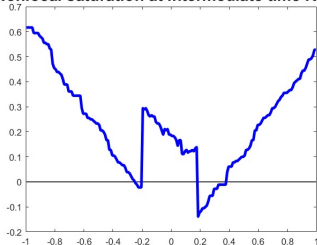
$$K = 0.6, w(r) = (R - r)/R$$



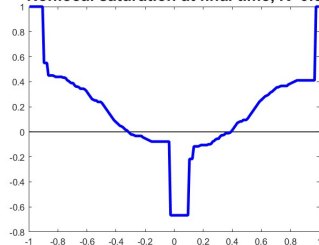


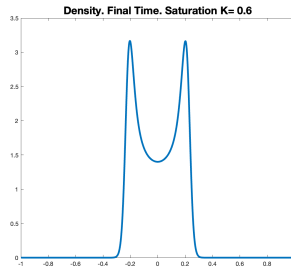
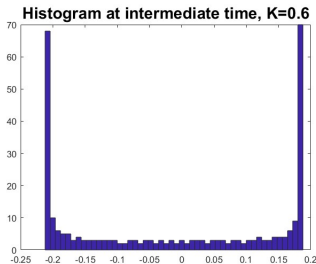
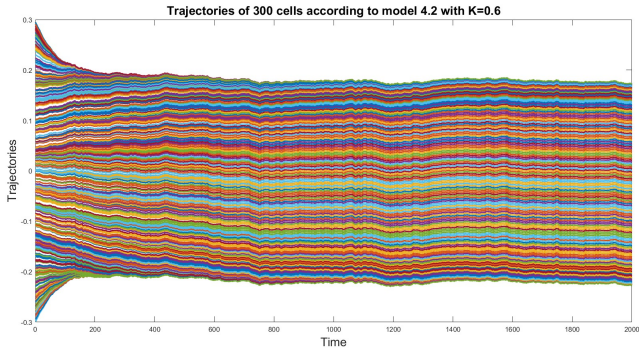


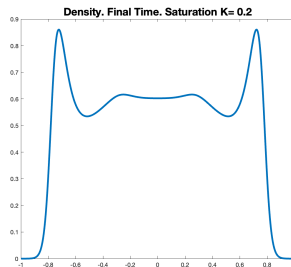
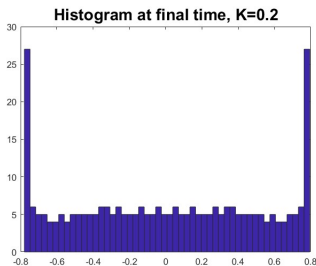
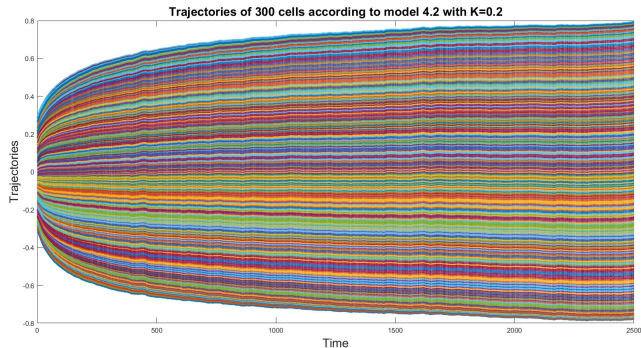
Nonlocal saturation at intermediate time K=0.6



Nonlocal saturation at final time, K=0.6







Present and future goals

- **Consider a long range diffusion of nonlocal nature:** Let $J(x, y)$ represent the probability of jumping from position x to y and $u(x, t)$ the density of cells as usual. If there is no other mechanisms than diffusion, the density evolution follows

$$u_t(x, t) = \int_{\text{supp}(J(x, \cdot))} J(x, y) u(y, t) dy - \int_{\text{supp}(J(\cdot, x))} J(y, x) dy u(x, t).$$

- **Scaling of a long range process with a differential process:** trajectories no longer differential.
 - **Protected/Preferred zones.**
 - Response to **external cues, rugosity of the surface, magnetic fields**, etc
- Analysis of the **sorting of two populations of cells** with the different models.
- Modelling **systems of Cell-Chemoattractant Substance.**

Thank you !!! Happy sixties for Kisko and Manolo !!!!

